

Microstructure evolution during warm working of Ti–6Al–4V with a colony- α microstructure

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Abstract

A high-resolution electron backscatter diffraction technique was employed to investigate microstructure evolution during warm working of Ti–6Al–4V with a colony- α microstructure. Particular emphasis was paid to the specific mechanisms governing this process. Microstructure development was found to be driven mainly by the geometrical requirements of the imposed strain and by the kinking of α lamellae. For the most part, the lamellar microstructure was surprisingly stable during straining, with limited globularization observed only in kinked α colonies. The kinking process was shown to be closely linked with the development of shear bands within the colonies. These observations suggest that changes in strain path may be beneficial in promoting globularization during warm working.

Keywords: Titanium alloys; Warm working; Microstructure; Electron backscatter diffraction

1. Introduction

The two-phase α/β titanium alloy Ti–6Al–4V is the backbone of the titanium industry especially for aerospace applications. This material is usually produced by an ingot-metallurgy route comprising a series of steps, each of which has specific microstructural goals. Following casting, ingots are typically hot worked and heat treated above the β transus (the temperature at which $\alpha + \beta \rightarrow \beta$) in order to eliminate casting inhomogeneities. Depending on the cooling rate following β processing, a variety of so-called transformed- β microstructures, ranging from those comprising coarse lamellar colonies to those consisting of very fine, acicular martensitic α , can be formed. Although such microstructures may promote resistance to fracture and high-temperature creep, a low-aspect ratio (so-called “globular”) microstructure is often desirable for applica-

tions requiring optimal strength or ductility. In this regard, lamellar microstructures are typically broken down to produce the desired globular ones via subtransus processing. This technological stage represents both the most difficult step from a workability standpoint as well as the most important with regard to final shaping and service.

Because of its great technological importance, the globularization process has been studied for many years. The majority of this research has focused on globularization at relatively high temperatures (800–950 °C) [1–7]. This prior work has established that there are generally two ways by which lamellae globularize during (and following) hot deformation [1,2]. Both involve the formation of boundaries across individual α lamellae and the subsequent diffusional penetration of the β phase along these boundaries, finally leading to break up of the lamellar α microstructure. In one case, the α/α grain boundaries originate from a recrystallization process, whereas in the second they are formed by deformation-induced intense, localized shear across the α lamellae.

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Recent work has shown that globularization at lower temperatures can lead to the formation of a finer $\alpha + \beta$ microstructure compared to that achieved at higher temperatures [8,9]. In fact, a submicrocrystalline structure has been produced by warm working at a temperature of ~ 550 °C [9]. Such refined microstructures may impart significantly higher strength at service temperatures and noticeably improved ductility at higher temperatures. Microstructure refinement in Ti–6Al–4V is therefore of considerable commercial interest. In addition, the physical basis of the globularization process at low temperatures is not clear, and the phenomenon is thus of substantial scientific interest as well.

Limited observations of Ti–6Al–4V have suggested several important features of the evolution of microstructure during deformation at low temperatures (550–700 °C) [8–10]. First, the majority of the α laths tend to become aligned in the principal direction of material flow, whereas some α colonies become kinked [10]. On a finer scale, the α laths become subdivided with increasing strain, leading to the formation of rows of relatively low-aspect ratio α particles/grains (or subgrains); however, details of this process are not clear [10]. Nevertheless, it has been demonstrated that globularization kinetics are significantly influenced by strain path. For a fixed strain path (e.g. extrusion), the lamellae microstructure is surprisingly stable, constituting ~ 60 vol.% of the material after true strains as large as 1.8 [10]. On the other hand, a change in strain path (e.g. via “abc” deformation) promotes the formation of a relatively homogeneous globular microstructure after true strains of ~ 3 [9].

Prior research has provided a broad insight into the globularization process at relatively low (warm-working) temperatures. However, these efforts have concentrated primarily on *final* microstructure with relatively little attention on the detailed mechanism of how microstructure evolves with deformation. The object of the present work, therefore, was to delineate the details of the mechanism by which globularization occurs during warm working of α/β titanium alloys with a lamellar preform microstructure. For this purpose, scanning electron microscopy (SEM) and electron backscatter diffraction (EBSD) were performed on Ti–6Al–4V samples deformed to different strain levels.

2. Materials and methods

The material used in the present investigation was commercial-grade Ti–6Al–4V whose measured composition (in wt.%) was 6.3 Al, 4.1 V, 0.18 Fe, 0.03 Si, 0.002 Zr, 0.01 C, 0.18 O and 0.01 N, balance Ti. The β transus for this material was approximately 995 °C. The material was supplied as 18 mm diameter hot-rolled bar with a duplex microstructure consisting of primary α in a transformed- β matrix. The as-received bar was heat treated using a cycle comprising 955 °C/15 min + 1010 °C/15 min + furnace cool to 800 °C/soak for 20 min + controlled furnace cool (at a rate of

2 °C min⁻¹) to 600 °C/soak for 2 h + water quench. By this means a colony- α preform microstructure was developed.

Cylindrical specimens 12 mm high and 8 mm in diameter were machined from the heat-treated bar. The specimens were uniaxially compressed to a height reduction of 25%, 50% or 70% at 600 °C and an initial strain rate of 10⁻³ s⁻¹ using a Schenk universal testing machine. After deformation, all specimens were water quenched to preserve the as-deformed microstructure. A sample was also heated to the deformation temperature, given an identical soak, and water quenched to determine the undeformed microstructure. Following deformation, each specimen was sectioned axially and prepared for SEM and EBSD. For SEM observations, the specimens were mechanically polished, electropolished in an 80%:20% acetic anhydride/perchloric acid solution at ~ 15 °C using a DC power supply operated at 70 V, and etched with Kroll's reagent (95% H₂O, 3% HNO₃, 2% HF). For EBSD, samples were mechanically polished followed by electropolishing in a solution of 5% perchloric acid, 35% butanol and 60% methanol at approximately -35 °C with an applied potential of 40 V.

All microstructure observations focused on the central (most highly deformed) region of the compression specimens; all SEM micrographs and EBSD scans were taken such that the compression direction was vertical. SEM studies were carried out using a JEOL JSM-6400. High-resolution EBSD analysis was conducted in a Hitachi S-4300SE FEG SEM equipped with TSL OIMTM software. Orientation mapping involved automated beam control using a triangular scanning grid; step sizes were 0.1–1 μ m. For each diffraction pattern, nine Kikuchi bands were used for indexing to minimize errors. In this manner, EBSD maps comprising $\sim 400,000$ points were obtained. The *average* confidence index (CI) for each EBSD map ranged from 0.66 to 0.12. By comparison, experiments on face-centered cubic (fcc) materials have shown that the fraction of correctly indexed patterns with CIs greater than 0.1 is 95% [11]. To obtain a more reliable picture of microstructure evolution, small grains comprising three or fewer pixels were automatically removed from the maps using the grain-dilation function in the TSL software. To eliminate spurious boundaries caused by orientation noise, a lower-limit boundary misorientation cutoff of 2° was used. All misorientation angles quoted below are relative to the rotation axis with the minimum misorientation. A misorientation of 15° was used to separate low-angle boundaries (LABs) from high-angle boundaries (HABs).

3. Results and discussion

3.1. General observations of microstructure evolution

General trends of microstructure evolution during warm deformation are summarized in Figs. 1 and 2. The effect of strain on microstructural morphology is illustrated by the SEM micrographs in Fig. 1; here, the hexagonal close-

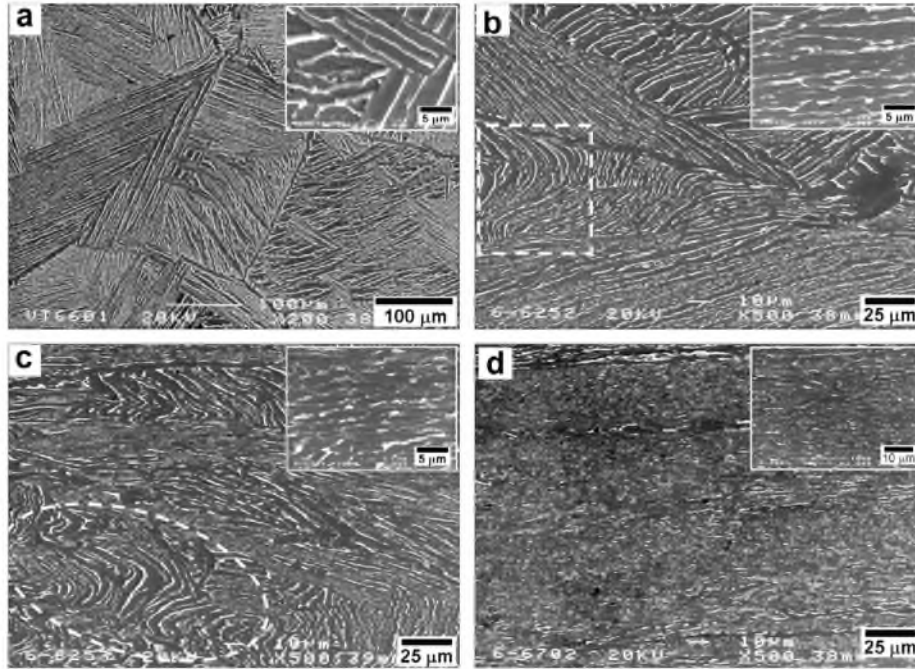


Fig. 1. SEM images of Ti-6Al-4V microstructure (a) prior to deformation, and after height reductions of (b) 25%, (c) 50% and (d) 70%. A high-magnification image of the microstructure is shown in the top right corner of each figure.

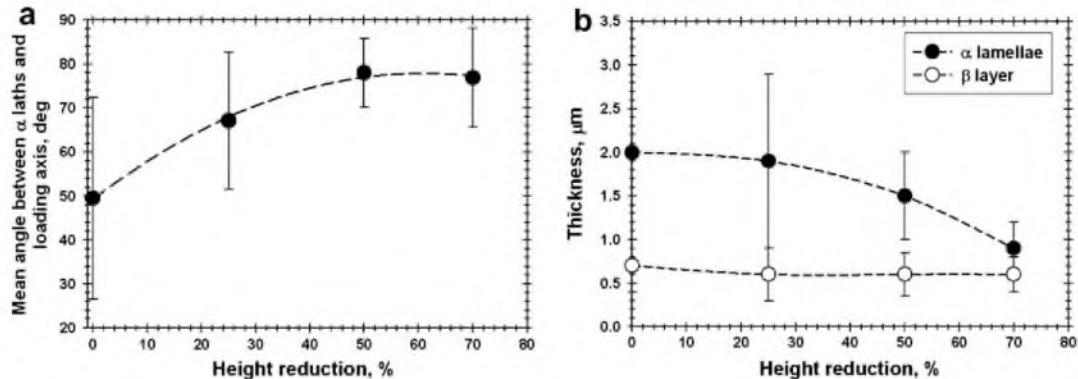


Fig. 2. Effect of deformation on (a) reorientation of α lamellae and (b) the thickness of α and β lamellae. Error bars show the standard deviations for the measurements.

packed (hcp) α phase is dark, and the body-centered cubic (bcc) β phase is the lighter constituent. The material quenched immediately prior to deformation had a "classical" transformed colony- α microstructure consisting of three microconstituents: lamellar/basketweave α colonies with 2 μm thick α platelets, layers of α at the prior- β grain boundaries, and thin ($\sim 0.6 \mu\text{m}$) β layers separating the α lamellae (Fig. 1a). The β volume fraction was measured to be $\sim 18 \pm 1\%$.

Microstructure evolution as a function of strain was found to be significantly influenced by the orientation of specific α colonies with respect to loading axis. If the plane normal to the α lamellae was nearly perpendicular to the compression axis, the α lamellae became kinked (e.g. circled areas in Fig. 1b and c). Otherwise, the α lamellae rotated with strain, with the α lamellae eventually becoming

aligned perpendicular to the loading axis (Fig. 2a). Such alignment of the α lamellae gradually eliminated the morphological differences between neighboring α colonies and resulted in the formation of a relatively homogeneous microstructure (Fig. 1c and d). This behavior is well documented for Ti-6Al-4V deformed at relatively high temperatures [4,7]. The primary difference for the present case involving warm working was that the β layers generally did not break up and were continuous in nature even at large strains (e.g. high-magnification inserts in Figs. 1b – d). These observations indicate that the globularization process is noticeably retarded. By comparison, dynamic globularization during hot working is typically completed at a height reduction of $\sim 80\%$ [4].

The effect of strain on the average thickness of the α and β lamellae is shown in Fig. 2b. The measured thickness of the α

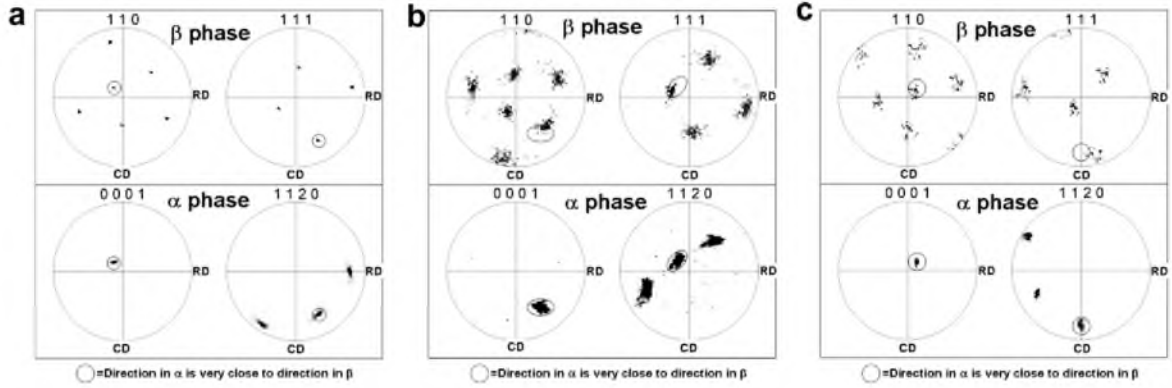


Fig. 3. Orientation relationship between the α and β phases in Ti-6Al-4V (a) prior to deformation, and after height reductions of (b) 25% and (c) 50%. CD and RD denote the compression axis and radial directions, respectively.

lamellae appeared to decrease with strain. It should be pointed out, however, that the measurements in Fig. 2b may not represent the true α platelet thickness because of stereological effects. In particular, if the plane of the platelets is not normal to the sectioning plane, the actual thickness of the α lamellae would be less than that measured on micrographs. This effect is greatest at low strains at which the α platelets lie at arbitrary angles to the sectioning plane. In such cases, the true α platelet thickness is approximately (1/1.5) times the measured thickness [12]. By contrast, at large strains, the α platelets rotate to an orientation close to being perpendicular to the compression axis (Fig. 2a) and the sectioning plane. In such cases, an axial section might indeed give a reasonable estimate of the platelet thickness. In view of these considerations, the actual rate of the α lamellae compression is thus less than the one shown in Fig. 2b.

The results in Fig. 2b also reveal that the average thickness of the β laths appeared to undergo little change with strain. This observation appears surprising because the β phase is generally believed to be softer than the α phase, at least at hot working temperatures. The deformation behavior of the β phase in Ti-6Al-4V thus deserves further investigation.

3.2. Burgers orientation relationship between the α and β phases

An important feature of the colony- α microstructure in α/β titanium alloys such as Ti-6Al-4V is the Burgers orientation relationship (OR) between the α and β phases in which the respective close-packed planes ($\{0002\}_\alpha/\{110\}_\beta$) and close-packed directions ($\langle 11\bar{2}0 \rangle_\alpha/\langle 111 \rangle_\beta$) are parallel. This relationship may provide easy slip transmission across α/β interfaces such that an individual α colony behaves as a single grain [13,14]. In the present work, EBSD provided insight into the local crystallographic relationship between the α and β phases and how it evolved with strain. To this end, orientation data for both phases were derived from high-resolution EBSD maps covering single α colonies; these data were used to construct 110 and 111 pole figures for the

β phase and 0002 and 11 $\bar{2}0$ pole figures for the α phase (Fig. 3).

In the undeformed material, the (0001) plane and one of $\langle 11\bar{2}0 \rangle$ directions of the α phase were nearly parallel to the one of the $\{110\}$ planes and one of the $\langle 111 \rangle$ directions in this plane of the β phase, respectively (Fig. 3a). Thus, the α and β phases were indeed related by the Burgers OR. Warm plastic deformation increased the orientation spread substantially in both phases and caused significant crystallographic rotations, leading to deviations from the exact Burgers OR at 25% height reduction (Fig. 3b) or elimination of it at 50% reduction (Fig. 3c).

Because the β layers were very thin ($\sim 0.6 \mu\text{m}$) and usually gave rise to diffuse Kikuchi patterns, it was difficult to ascertain whether the examples in Fig. 3 were truly representative. To provide evidence of the elimination of the OR which was more statistically reliable, the boundaries between α colonies within given β grains were examined. In particular, due to the crystallography of the α and β phases, there are 12 possible α orientations (each satisfying the Burgers OR) that can transform from a single parent β grain, i.e. there are six $\{110\}$ planes in a given β grain, each containing two $\langle 111 \rangle$ directions, on which an α variant/colony can form. The possible misorientations between α variants/colonies inherited from the same parent β grain are summarized in Table 1. It is seen that these misorientations are typically very specific, and their unique nature is thus indicative of the Burgers OR.

Fig. 4 shows EBSD image quality (IQ) maps in which the boundaries between the α variants for both the starting and the deformed material have been highlighted. The IQ map quantifies the sharpness of the Kikuchi pattern at each point and is similar to an SEM image with the difference that the β phase appears dark and the α phase light. Fig. 4a shows that the colony boundaries in the initial material formed a continuous network. Using a 5° tolerance criterion, the boundaries between the α variants accounted for $\sim 70\%$ of the total boundary area in the α phase, thus confirming the Burgers OR between the α and β phases (Fig. 3a). After 25% height reduction, the boundary network became discontinuous (Fig. 4b). This

Table 1
Misorientation between α variants inherited from the same parent β grain [15,16].

Misorientation (angle-axis pair)	Symbol	Location of misorientation axis on stereographic triangle
$60^\circ < 11\bar{2}0 >$	●	
$60.8^\circ < \bar{1}.377; \bar{1}; 2.377; 0.359 >$	○	
$63.3^\circ < \bar{1}0; 5; 5; \bar{3} >$	■	
$90^\circ < 1; 2.38; 1.38; 0 >$	□	
$10.5^\circ < 0001 >$	▲	

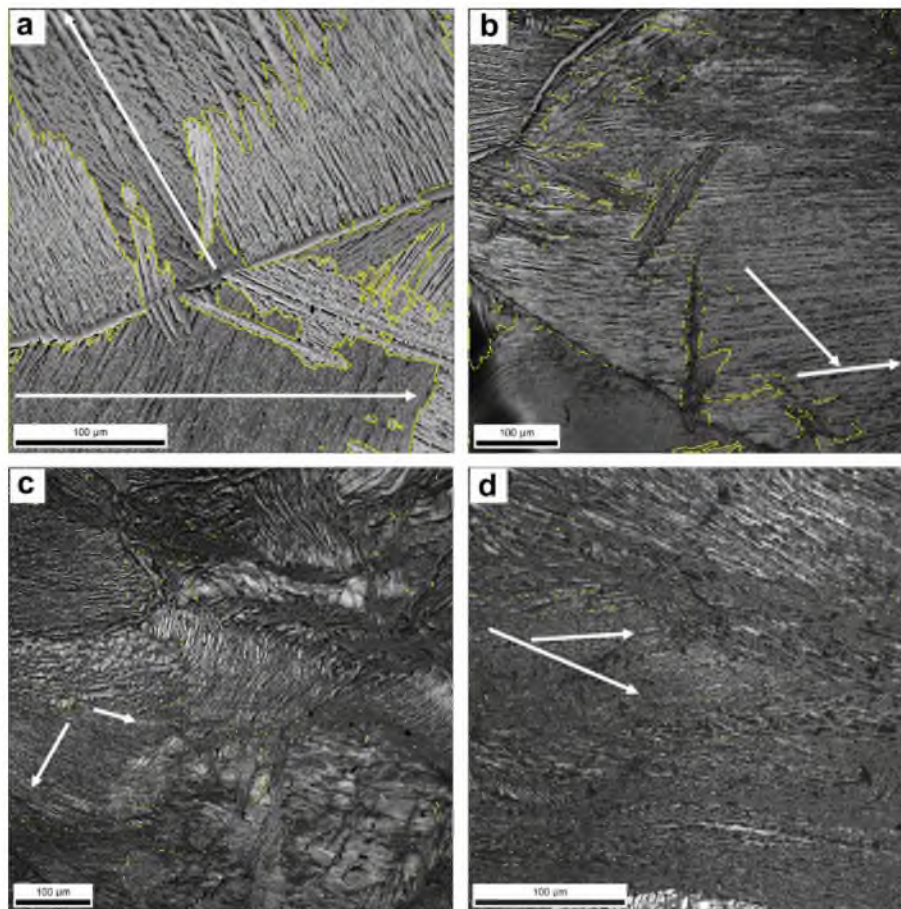


Fig. 4. EBSD image quality maps with highlighted boundaries linked to the Burgers orientation relationship in Ti-6Al-4V (a) prior to deformation, and after height reductions of (b) 25%, (c) 50% and (d) 70%. The arrows indicate the locations of misorientation line scans shown in Figs. 5 and 6.

observation suggests that adjacent colonies may have rotated relative to one another, producing localized flow in the boundary between them and hence loss of the OR. This process is similar to the evolution of twin misorientations in pure titanium during straining [17] and may be explained by local deformation-induced crystallographic rotation of each α colony away from its initial orientation inherited during the β -to- α phase transformation. This result supports the conclusion that the breakdown of the Burgers OR between the α and β phases began at reduc-

tions of the order of 25% (Fig. 3b). In the material deformed to a height reduction of 50–70%, the inter-colony boundaries have disappeared almost completely, and only a few local boundary segments were noted (Fig. 4c and d).

3.3. Orientation spread within α colonies

The breakdown of the Burgers OR was associated with significant crystallographic rotations of the α and β phases within the α colonies. In order to provide a deeper insight

into this phenomenon, orientation variations transverse and parallel to α laths were measured in EBSD maps such as those shown in Fig. 4; typical examples of the cumulative misorientation are given in Figs. 5 and 6, respectively.

In the undeformed material (Figs. 5 and 6a), orientation variations in both the transverse and parallel directions were generally within the angular resolution of EBSD ($\sim 2^\circ$). Thus, the α laths within a given colony had essentially the same crystallographic orientation, and the orientation spread was very small. These observations were as expected and may be readily explained in terms of the Burgers OR governing the β -to- α phase transformation (Fig. 3a).

Plastic deformation gave rise to a significant orientation spread within each α colony as well as within individual α laths. The misorientation profile in the direction *transverse* to the α laths (Figs. 5b–d) showed alternating orientation changes with a “wavelength” of ~ 5 – $10\ \mu\text{m}$, a length that is comparable to the initial α lath thickness of $\sim 2\ \mu\text{m}$. This suggests that groups of neighboring α laths rotated relative to each other, and these rotations tended to compensate each other. This trend is reminiscent of observations for deformed single-phase fcc metals (e.g. [18,19]). It is believed that this effect in the fcc materials is associated with the activation of different slip systems in adjacent cell blocks [20].

It is also important to note that neighboring α laths started to rotate relative to each other at a height reduction as low as 25% (Fig. 5b). At this strain level, the Burgers OR

between the phases was still somewhat evident (Figs. 3 and 4b), and hence the α colonies should have behaved as a single grain [13,14]. However, this was not the case, and thus the Burgers OR itself does not ensure homogeneous deformation within the α colonies during warm working.

EBSD profiles measured *along the length* of α laths revealed broadly increasing misorientations with distance (Fig. 6b–d). The misorientation profiles were typically relatively smooth, suggesting that the orientation change was continuous in nature rather than one associated with the formation of HABs. Furthermore, the magnitude of the orientation gradient was found to increase with strain, reaching $\sim 5^\circ\ \mu\text{m}^{-1}$ at a height reduction of 70% (Fig. 6d). The continuous character of the orientation spread assumes a curvature of the crystalline lattice (i.e. disclination) with the associated (internal) stress $\sigma \sim G\omega$, where G is the shear modulus and ω is the cumulative rotation angle (Frank vector) (e.g. [21]). The internal stress appears to increase abruptly in a strain range between 25% and 50% height reduction (Fig. 6b and c). At this strain level, the Burgers OR is destroyed (Figs. 3b, c and 4b, c). Therefore, these two phenomena may be related.

3.4. Grain boundaries in the α phase

In order to provide fundamental insight into the globularization process, it is necessary to study the development of grain boundaries in the α phase. Representative EBSD

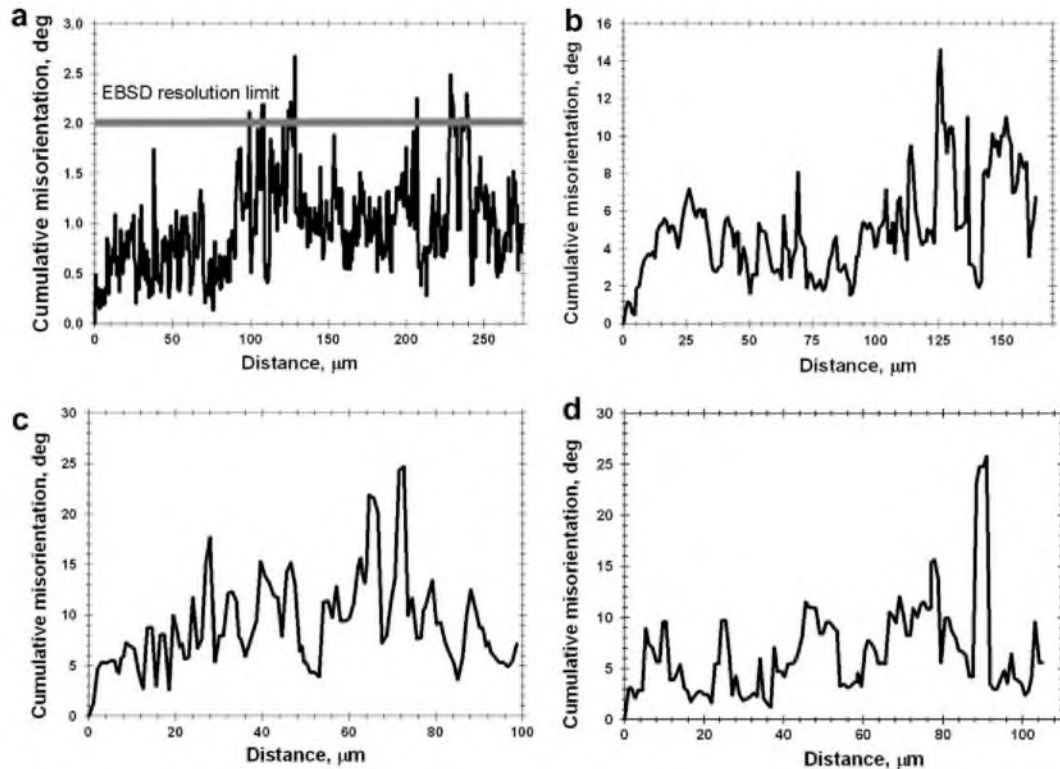


Fig. 5. Typical line scans from EBSD maps in Fig. 4 showing the variation of orientation over a long distance transverse to the α laths in Ti-6Al-4V (a) prior to deformation and after height reductions of (b) 25%, (c) 50% and (d) 70%.

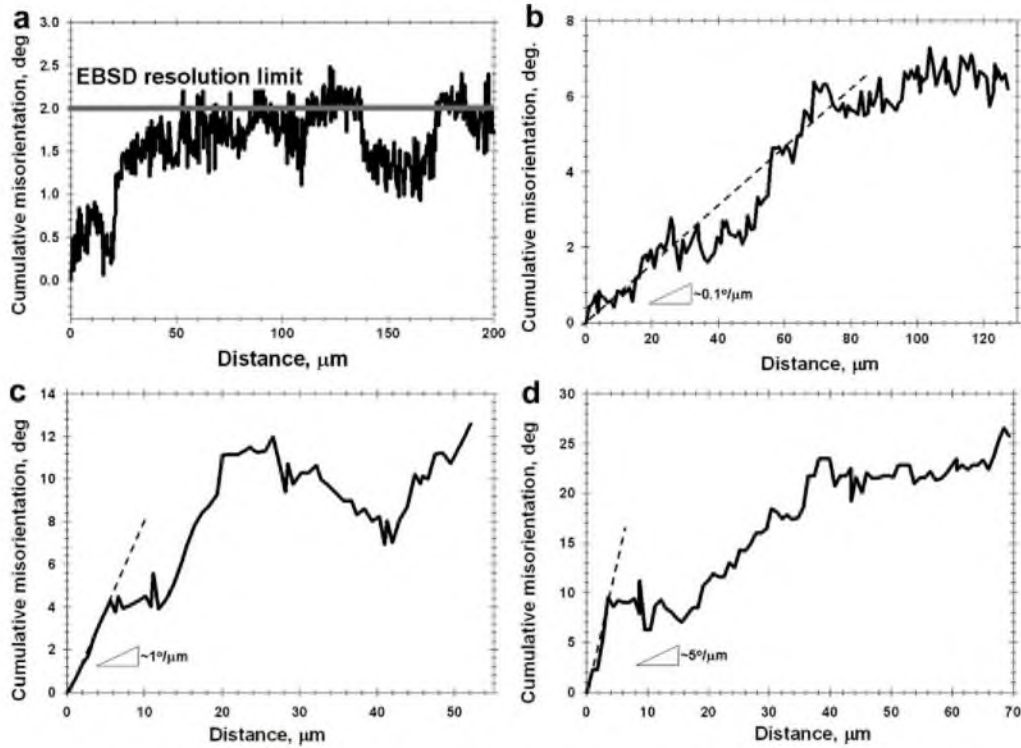


Fig. 6. Typical line scans from EBSD maps in Fig. 4 showing the variation in orientation over a long distance along the α laths in Ti-6Al-4V (a) prior to deformation, and after height reductions of (b) 25%, (c) 50% and (d) 70%.

maps showing the arrangement of LABs and HABs in the starting and deformed materials are given in Fig. 7. In the maps, all possible α grain boundaries are shown, including those between α variants/colonies as well as the boundaries within/between the α laths; LABs and HABs are depicted as red and black lines, respectively. Misorientation data derived from large EBSD maps (including several thousands of α laths each) are summarized in Fig. 8.

In the *undeformed* material, most of the boundaries were high angle in character; only sporadic LABs were found within some α colonies (Fig. 7a). The misorientation distribution was dominated by 10° , 60° and 90° boundaries whose rotation axes were clustered near the $[0001]$ and $\langle 2\bar{1}\bar{1}0 \rangle$ poles (Fig. 8a). A comparison of the measured misorientations (Fig. 8a) and those predicted for the Burgers OR (Table 1) showed very close agreement. As mentioned in Section 3.2, the fraction of the boundaries between the α variants in the undeformed material was 70%, thus indicating that these boundaries originated primarily from the β -to- α phase transformation involving the Burgers OR. The remaining 30% of the α grain boundaries included some LABs and boundaries between parent β grains.

Plastic deformation dramatically increased the fraction of LABs (Fig. 7b–d) skewing the misorientation-angle distribution towards low-angle rotations (Fig. 8b–d). After a 25% height reduction, the distribution still exhibited several misorientation peaks associated with the Burgers OR (Fig. 8b), but larger strains led to the unambiguous dominance of maxima at low misorientation angles (Figs. 8c and

d). Furthermore, with increasing strain, the misorientation across some boundary segments began to reach and exceed 15° . This trend was most obvious for the sample subjected to the 70% height reduction (Fig. 8d). The general deformation of the α colonies brought about by the imposed compressive strain also contributed to an increase in the fraction of HABs (Fig. 8b–d).

It should be emphasized, however, that the LABs in the EBSD maps were not located randomly within the microstructure, but were aligned primarily with the longitudinal direction separating neighboring α laths (Fig. 7b–d). In reality, neighboring α laths within an α colony have no common grain boundary in the actual microstructure because they are usually separated by a continuous layer of β , as shown in the high-magnification inserts in Fig. 1b–d. This observation may therefore be explained in terms of the low resolution of the β phase in EBSD, as mentioned above. In other words, the β phase is frequently missed during EBSD or “cleaned” during data processing. It is likely therefore that the longitudinal boundaries seen in the EBSD maps are somewhat artificial and simply reflect the orientation difference between neighboring α laths, as shown in Fig. 5b–d.

It is noteworthy that rotation axes of the $85\text{--}90^\circ$ boundaries tend to cluster near $\langle 2\bar{1}\bar{1}0 \rangle$ and $\langle 10\bar{1}0 \rangle$ poles (not shown). This may be associated with $\{10\bar{1}2\}$ and $\{11\bar{2}3\}$ twinning [16]. An example of a possible $\{11\bar{2}3\}$ twin is indicated by the arrow in Fig. 7c. It should be noted, nevertheless, that twin identification requires determination of not only the grain-boundary misorientation but also the twin boundary plane. Misorientation data

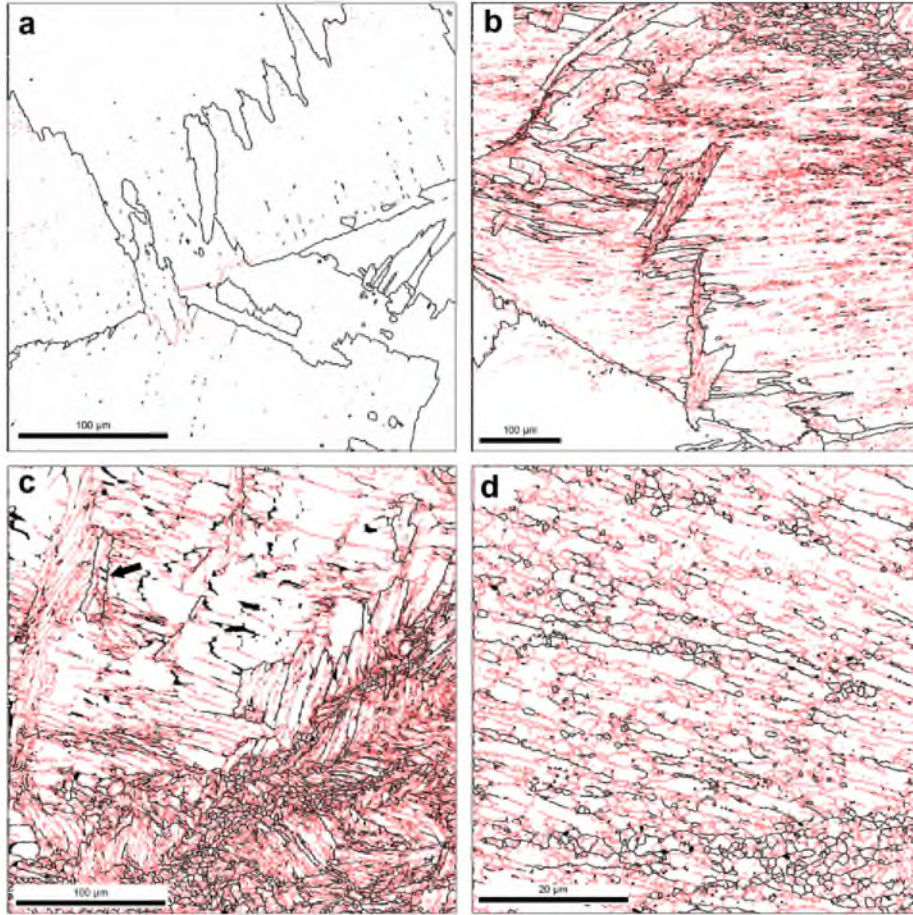


Fig. 7. EBSD results showing the low- and high-angle grain boundaries in the α phase (depicted as red and black lines, respectively) (a) prior to deformation, and after height reductions of (b) 25%, (c) 50% and (d) 70%. Note that the small amount of β phase appears black in the figures.

are routinely provided by EBSD but grain boundary planes cannot be measured due to the two-dimensional nature of EBSD maps. Hence, the possible twinning during warm working requires additional verification.

3.5. Slip in the α phase

The observation that the clustering near the $[0001]$ pole in misorientation-axis distributions (Fig. 8b–d) is mainly produced by boundaries with misorientations in the range of $5\text{--}15^\circ$ is of particular importance. Because the formation of these misorientations was mainly associated with the crystallographic rotation of neighboring α laths, this finding provides insight into the nature of deformation and slip in the colony microstructure. It is generally accepted that the activity of slip system p rotates the crystalline lattice around the axis $w_p \parallel (b \times n)_p$, in which b denotes the Burgers vector magnitude, and n is the slip plane normal. Among the possible slip modes in α titanium (prism, basal and pyramidal), *only* prism slip provides a lattice rotation about the $[0001]$ direction. Thus, the proximity of the rotation axes to the $[0001]$ pole indirectly suggests preferential activation of prism slip, which is commonly accepted to be the dominant deformation mode in α titanium [22]. How-

ever, prism slip can provide only two independent slip systems: $(10\bar{1}0)[1210]$ and $(1\bar{1}00)[11\bar{2}0]$. By using the rotation-axis equation, therefore, it may be shown that these slip systems tend to rotate the crystalline lattice in opposite directions. Hence, it is reasonable to hypothesize that the alternating nature of the misorientations between neighboring α laths (Fig. 5b–d) is indeed associated with the activation of different slip systems; this behavior is analogous to observations for fcc metals [18–20]. It is likely that material flow in each α lamella arises mainly from one dominant slip system (of the two independent systems) which rotates neighboring α laths in mutually opposite directions.

The question thus arises as to why different slip systems operate in the neighboring α laths despite their very close initial crystallographic orientation (Fig. 5a)? It is well accepted that five independent slip systems are generally required to accommodate any arbitrary imposed strain. On the other hand, flow stress and strain hardening are known to increase with the number of the operating slip systems (because the number of intersecting dislocations increases), and thus there is an energetic benefit in reducing their number. In this regard, Hansen et al. [20] developed a model based on the subdivision of a deformed grain into

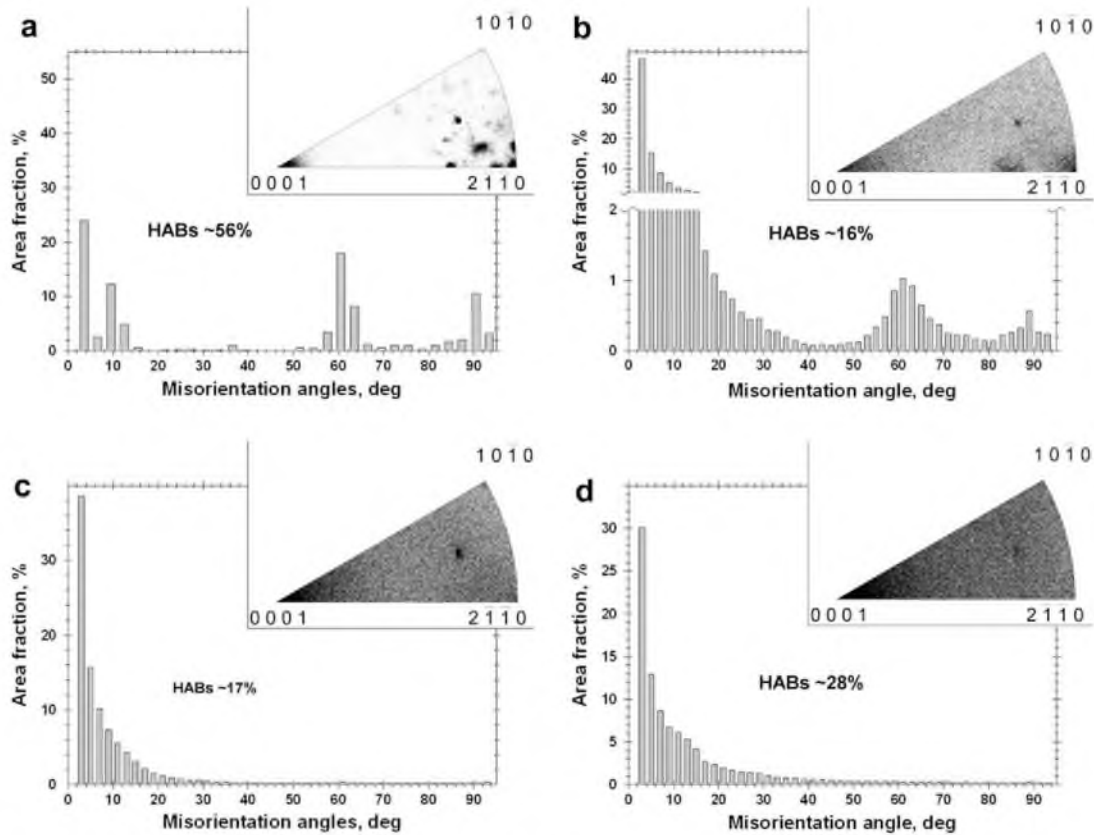


Fig. 8. Misorientation distribution in the α phase (a) prior to deformation, and after height reductions of (b) 25%, (c) 50% and (d) 70%. The corresponding misorientation-axis distribution is shown in the top right corner of each figure.

two-dimensional cell blocks, each of which deforms on fewer than five slip systems; however, the Mises requirement is globally satisfied on the grain scale by the overall deformation behavior of the cell blocks within a given grain [20]. In the present case, a similar situation appears to apply. Because of the very specific morphology of the transformed- β microstructure, the α colonies are already separated into two-dimensional α lamellae. Additionally, the sample-scale strain hardening during warm working was close to zero for deformations between 25% and 70% height reduction (Fig. 9). It is likely therefore that the activation of different slip systems in the neighboring α laths is associated with limited strain hardening.

The dominance of only *one* alternating slip system (of the two independent systems) in neighboring α lamellae may efficiently rotate the α laths in the mutually opposite directions. On the other hand, single slip would strongly inhibit the formation of dislocation boundaries within the α lamellae as the activation of two or three slip systems is usually required to produce stable dislocation boundaries; single slip may only give rise to unstable walls of edge dislocations. Close inspection of Fig. 8b–d confirms this hypothesis. The majority of LABs are indeed oriented along the α lamellae (i.e. they are artificial, as explained above), and only a small fraction of the (real) LABs lie in the transverse direction. Thus, slip “sharing” between

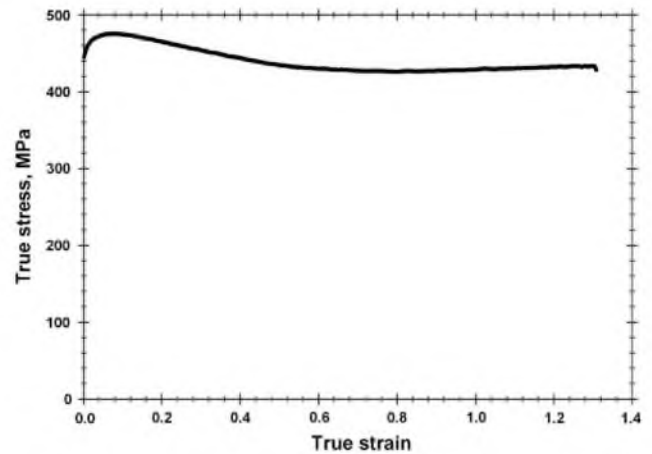


Fig. 9. Typical true stress–strain curve for warm working of Ti-6Al-4V with colony- α microstructure.

neighboring α lamellae appears to inhibit the globularization of the α phase.

Irrespective of the number of slip systems operating within each α lamella, strain accommodation on an α -colony level is also required. Taking into account that only two independent prism slip systems are available, the question arises as to how many slip systems are needed to satisfy this strain-compatibility requirement? Kocks and

Canova [23] have shown that, under some specific conditions, strain compatibility can be achieved with far fewer than five slip systems. For example, during the compression of flat grains having *random* orientations, the activation of the five systems is only needed at the periphery of each grain, whereas the bulk material may deform on only three independent slip systems [23]. In the present work, the original α colonies are initially equiaxed (Fig. 1a). However, the number of necessary slip systems would be expected to decrease by deformations in the range of ~ 50 – 70% height reduction when the colonies become pancaked. Bearing in mind that plastic deformation in the α phase is usually associated with the formation of some texture (i.e. orientation of the neighboring α colonies is expected not to be random), it may be hypothesized that the number of necessary slip systems decreases to the two that are available in this strain range. It should also be borne in mind that the deformation is not identical from one colony to another, i.e. some colonies form kinks and others do not. This strain nonuniformity may thus enable the activation of fewer than five slip systems within each colony.

Strain compatibility requirements appear to be most severe at the peripheral parts of the α laths because strain accommodation between the neighboring α colonies is needed. The potential lack of the available slip systems in these regions may give rise to internal stress and an associated orientation spread in the longitudinal direction of the α lamellae (Fig. 6b–d). On the other hand, the activation of the polyslip at the edges of α lamellae may stimulate globularization in these regions. An example may be seen in the bottom right corner of Fig. 7d. This effect has also been observed during the hot working of $\alpha + \beta$ titanium alloys [24].

3.6. Kinking of α colonies

The grain-boundary maps in Fig. 7b–d showed that the α laths contained only a limited number of transverse boundaries, and most of these were low angle in nature. This finding underscores the observation that the rate of dynamic globularization of the α phase at warm-working temperatures was much slower than in the hot-working regime. In fact, only a

small fraction of globularized microstructure was found in the deformed material (Figs. 7c and d).

The globularized regions were typically found in areas in which the colonies had formed kinks. Hence, high-resolution EBSD was applied to gain insight into the details of the kinking process. To obtain maximum information, each stage of kinking was characterized by both an IQ map (Fig. 10) and an inverse-pole figure (Fig. 11). As noted earlier, IQ maps are similar to conventional SEM images (with the difference that the β phase appears dark) and thus can serve as a reference in the present discussion. Inverse-pole-figure maps, on the other hand, convey detailed crystallographic information and thus are useful for in-depth insight into microstructure development. In the inverse-pole-figure maps presented here, individual α grains are colored according to their crystallographic orientation relative to the radial direction: the β phase appears black; and LABs and HABs in the α phase are depicted with white and black lines, respectively.

The early stage of the kinking process is shown in Figs. 10a and 11a. At this stage, the β layers had a zigzag appearance (e.g. circled area in Fig. 10a). A close inspection (e.g. circled area in Fig. 11a) showed that such zigzag patterns resulted from shear bands extending over long distances within an α colony; these shear bands crossed both the α laths and β layers. The variation of orientation color indicated that material on either side of the shear bands experienced a crystallographic rotation. The bands were not completely delineated by LABs; thus, some boundary segments probably had misorientations lower than the resolution limit of EBSD.

As the kinks became sharper (Fig. 10b), the shear bands became more pronounced and were more fully outlined by LABs (Fig. 11b); moreover, some boundary segments had accumulated a high-angle misorientation (e.g. region circled in Fig. 11b). The progressive development of the shear bands continued during the kinking process, and, in severely kinked α colonies (Fig. 10c), the bands were completely delineated by high-angle grain boundaries (Fig. 11c). The development of the shear bands within a kinked α colony also led to significant shearing (and break up) of the β layers (Fig. 10c) and produced a low-aspect ratio α phase grain/subgrain microstructure (Fig. 11c).

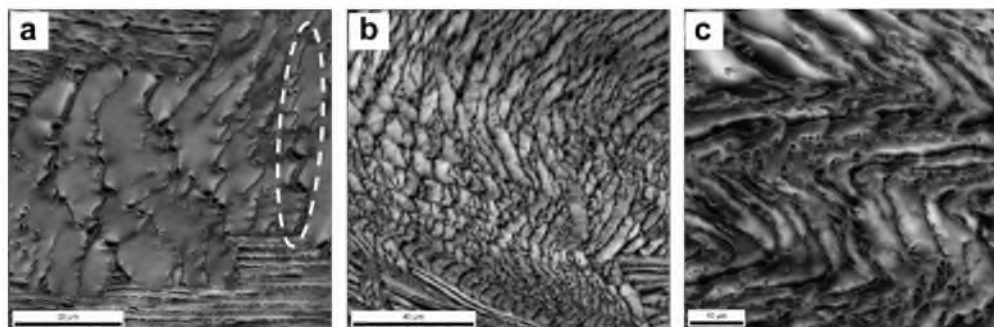


Fig. 10. EBSD image quality maps illustrating the progressive development of the kinking process in α lamellae. The compression direction is vertical.

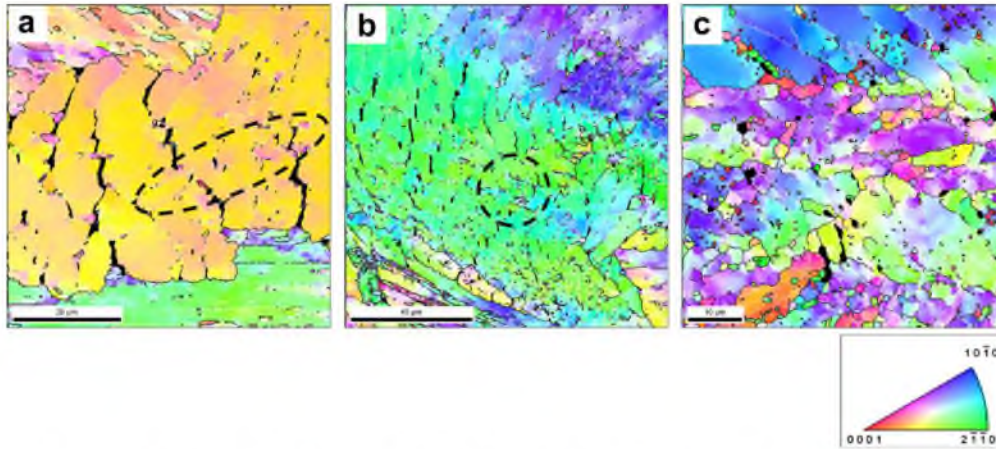


Fig. 11. EBSD inverse-pole-figure maps for radial direction illustrating the progressive development of the kinking process in α lamellae. In the maps, LABs and HABs in the α phase are depicted as white and black lines, respectively, and the β phase appears black. The compression direction is vertical.

Thus, the kinking process was closely linked with the development of shear bands within α colonies. Similar behavior has been noted for high-temperature deformation [6,24]. Shear banding therefore appears to be an intrinsic feature of kinking.

It is noteworthy that there was no apparent penetration of the β phase into the α phase even in severely kinked α colonies (Figs. 1c and 10c). This observation, in conjunction with the noticeable level of internal stress in the α phase (deduced from Fig. 6c–d), suggests that diffusional processes were not very active. Thus, globularization during warm working was sporadic in nature in contrast to the corresponding high-temperature behavior. It may thus be concluded to arise mainly from deformation with a negligible contribution from diffusional mechanisms.

3.7. Implications for efficient warm working

Extensive microstructure observations for Ti-6Al-4V have shown that a constant-strain path is the most effective method for globularizing of a colony- α microstructure at *hot-working temperatures* [25]. The present work, however, has demonstrated that the lamellar structure is very stable during warm deformation (Fig. 7b–d). Although the physical basis of this phenomenon is far from being completely understood, it appears that this effect may be associated with the limited number of slip systems operating within α laths due to slip sharing between neighboring α lamella. Nevertheless, a small amount of globularization is achieved through the kinking of the α laths whose plane normal is nearly perpendicular to the loading direction (Fig. 11c). Statistical measurements indicated that the remaining α laths tended to be reoriented in a position with the plane normal nearly parallel the compression direction after $\sim 50\%$ height reduction (Fig. 2a). Hence, a change in strain path comprising rotation of the workpiece at this level of deformation may enhance the rate of globularization during warm deformation.

4. Conclusions

The EBSD technique has been employed to investigate microstructure evolution during warm working of Ti-6Al-4V alloy with a colony- α microstructure. The main conclusions from this work are as follows.

1. Microstructure evolution is driven mainly by the geometric requirements of the imposed strain and by the kinking of the α lamellae. If the plane normal to the α lamellae was nearly perpendicular to the compression axis, the α laths are kinked. Otherwise, the α laths rotate with strain, with the plane normal tending to become aligned parallel to the loading axis, and their thickness is reduced somewhat.
2. During warm working, material flow in the α phase arises mainly from prism slip.
3. The crystallography of the transformed- β microstructure is governed by the Burgers OR; local crystallographic orientations of the α and β phases are related via the Burgers OR and the misorientation distribution in the α phase is dominated by the boundaries between the α variants. Plastic deformation destroys the Burgers OR between the phases and randomizes the boundaries between the α variants.
4. Warm plastic deformation is inhomogeneous within α colonies as well as within α laths. Groups of neighboring α laths are rotated with respect to each other and these rotations tend to compensate each other. It is hypothesized that this phenomenon is associated with the limited number of slip systems in α titanium which is insufficient to accommodate an arbitrary imposed strain. On the other hand, significant orientation gradients are developed along individual α laths during deformation, suggesting the development of substantial internal stresses.
5. The lamellar structure is surprisingly stable during warm deformation; limited globularization is found, primarily within kinked regions. The kinking process is closely linked with the development of shear bands.

6. In contrast to high-temperature deformation, a change of strain path may provide a promising method to promote globularization during warm working.

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